

## Werkcollege, Sterrenstelsels, Week 7

These are the assignments for week 7 of the course *Sterrenstelsels*.

Every week, one of the problems provides credit towards the final exam. If at least **five** of these problems are handed in and approved, one question on the final exam may be skipped. The hand-in assignment for this week is **Problem 7.3** below.

### 7.1 Chemical evolution with inflow

#### Adapted from Sparke & Gallagher, problem 4.10

Suppose that the inflow of metal-poor gas is proportional to the rate at which new stars form, so that

$$\delta M_s + \delta M_g = \nu \delta M_s \quad (7.1.1)$$

for some constant  $\nu > 0$ .

- Verify Eq. (7.1.1)
- Show that S&G Equation 4.14,

$$\delta Z = \delta \left( \frac{M_h}{M_g} \right) = \frac{p \delta M_s - Z [\delta M_s + \delta M_g]}{M_g} \quad (7.1.2)$$

becomes

$$\delta Z = \frac{p - \nu Z}{\nu - 1} \frac{\delta M_g}{M_g} \quad (7.1.3)$$

so that the metallicity of the gas is

$$Z = \frac{p}{\nu} \left[ 1 - \left( \frac{M_g(t)}{M_g(0)} \right)^{\nu/(1-\nu)} \right] \quad (7.1.4)$$

### 7.2 Two-body relaxation

- Chapter 3.2.2: Fill in the missing steps leading up to (3.50) and (3.51).  
You may find the following integral useful:

$$\int \frac{1}{(A + Bx^2)^{3/2}} dx = \frac{x}{A \sqrt{A + Bx^2}} \quad (7.2.1)$$

We saw that the two-body relaxation time scale can be written as

$$t_{\text{rel}} = \frac{V^3}{8\pi G^2 m^2 n \ln \Lambda} \quad (7.2.2)$$

for relative velocities  $V$ , particle mass  $m$  and particle density  $n$ . Being slightly more precise about the meaning of  $V$ , we rewrite this as

$$t_{\text{rel}} = \frac{\langle V^2 \rangle^{3/2}}{8\pi G^2 m^2 n \ln \Lambda} \quad (7.2.3)$$

For a system in *virial equilibrium*, the velocity dispersion  $\langle V^2 \rangle$ , mass  $M$  and radius  $R$  of the system are related as

$$M = \eta \frac{\langle V^2 \rangle R}{G} \quad (7.2.4)$$

The constant  $\eta$  depends on the radial structure of the system and how exactly the radius  $R$  is defined (note the similarity of this relation to the corresponding relation for a circular orbit).

- For a system of uniform density, total mass  $M$  and radius  $R$ , show that the two-body relaxation time scale for a virialized system can be expressed as

$$t_{\text{rel}} = \frac{M^{1/2} R^{3/2}}{6G^{1/2} \eta^{3/2} m \ln \Lambda} \quad (7.2.5)$$

- With  $M$  and  $R$  defined as above, the constant  $\eta = 5/3$ . Now calculate  $t_{\text{rel}}$  for the following systems (you can assume a mean stellar mass of  $m = 1M_{\odot}$  and  $\ln \Lambda = 10$  in all cases):
  - An open cluster of stars ( $M = 10^3 M_{\odot}$ ,  $R = 3$  pc)
  - A globular cluster ( $M = 10^6 M_{\odot}$ ,  $R = 10$  pc)
  - A giant elliptical galaxy ( $M = 10^{12} M_{\odot}$ ,  $R = 10$  kpc)

### 7.3 Spiral structure: Lindblad resonances

In the theory of spiral structure, the *global pattern speed*  $\Omega_{\text{GP}}$  is the overall pattern speed of the spiral density wave. In general, this may differ from the *local pattern speed* for a two-armed spiral pattern,  $\Omega_{\text{LP}}$ , that follows from the condition  $\Omega_{\text{LP}} = \Omega - \frac{1}{2}\kappa$ . However, where the two are equal, we talk about an *Inner Lindblad Resonance*.

A spiral galaxy is observed to have a flat rotation curve with circular speed  $v_c = 250 \text{ km s}^{-1}$  from near the centre of the galaxy to very large radii. The co-rotation radius of the spiral pattern is at  $R = 5 \text{ kpc}$ .

- Calculate the location of the Inner ( $m = 2$ ) Lindblad Resonance. To do this, you will need to know how to compute the epicyclic frequency from the rotation curve. This can be done using the Oort constants:

$$\kappa^2 = -4B(A - B) \tag{7.3.1}$$

for Oort constants  $A$  and  $B$  (see chapter 3.3 in S&G).

- Assuming that the mass distribution is dominated by a spherical dark matter halo, calculate the density profile  $\rho(R)$  of the halo.